

UBC Physics 102

Lecture 5 Version 2

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Electric field [Text: Sect. 21-6]

- **Definition: electric field**
- If force \mathbf{F} on test charge q then electric field \mathbf{E} is

$$\mathbf{E} = \frac{\mathbf{F}}{q}$$

- Force depends on charge q but \mathbf{E} is the same for *all test charges*.
- So electric field is more useful quantity to work with. Once you know \mathbf{E} can easily compute force on any test charge q via

$$\mathbf{F} = q\mathbf{E}.$$



Outline

- ▷ Electric field
- ▷ Conductors
- ▷ Continuous charge distributions
- ▷ End



Electric field, contd

- **Definition: Coulomb's law**
- Convenient to use electric field form of Coulomb's law.
- Gives field at any point due to charge Q ,



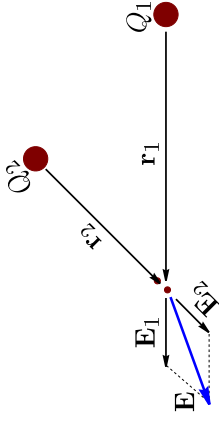
$$\mathbf{E} = k \frac{Q}{r^2} \hat{\mathbf{r}}.$$



Electric field, contd

• Discussion: Superposition principle

- If dealing with more than one charge, can just add up electric field due to each to calculate net electric field at a point,



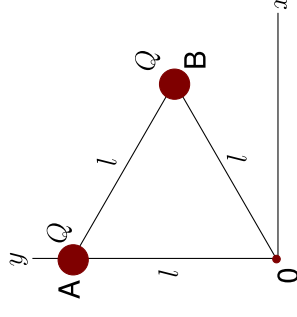
$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2 + \dots$$



Electric field, contd

• Example: Pr. 40

- Determine the electric field \mathbf{E} at the origin 0 due to the two charges at A and B.



Electric field, contd

• Solution: Pr. 40

- By the superposition principle

$$\mathbf{E} = \mathbf{E}_A + \mathbf{E}_B = E_A \hat{\mathbf{r}}_A + E_B \hat{\mathbf{r}}_B.$$
- From Coulomb's law the magnitudes of the electric fields are

$$E_A = E_B = \frac{kQ}{l^2}.$$
- Just need to find the directions. Direction from A to origin is

$$\hat{\mathbf{r}}_A = -\hat{\mathbf{j}}.$$



Electric field, contd

• Solution: Pr. 40, contd (Correction)

- B is at (x, y) where $y = \frac{1}{2}l$ and $x^2 + y^2 = l^2$ so

$$x = \sqrt{\frac{3}{4}}l.$$
- Direction from B to origin is

$$\hat{\mathbf{r}}_B = -\sqrt{\frac{3}{4}}\hat{\mathbf{i}} - \frac{1}{2}\hat{\mathbf{j}} = -\frac{\sqrt{3}}{2}\hat{\mathbf{i}} - \frac{1}{2}\hat{\mathbf{j}}.$$
- So net electric field at origin is

$$\begin{aligned} \mathbf{E} &= \frac{kQ}{l^2} (\hat{\mathbf{r}}_A + \hat{\mathbf{r}}_B) \\ &= \frac{kQ}{l^2} \left(-\frac{\sqrt{3}}{2}\hat{\mathbf{i}} - \frac{3}{2}\hat{\mathbf{j}} \right). \quad \square \end{aligned}$$



Conductors [Text: Sect. 21-9]

- **Interactive Quiz: PRS 05a**
- **Discussion: Conductors**
 - Conductors have free electrons.
 - Electrons move under force of electric field until the electric field is zero.
 - So electric field inside a conductor is always zero (after electrons have reached final position).
 - If a conductor has a net charge, it is always distributed on the surface, never in the interior the conductor.

- **Interactive Quiz: PRS 05b**

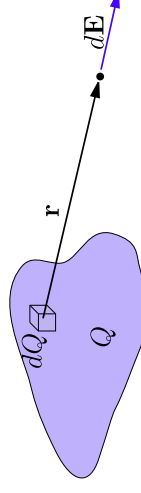


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Continuous charge distributions [Text: Sect. 21-7]

- **Discussion: Continuous charges**
 - If object too large to be treated as point charge, can still solve for electric field.
 - Divide object into small chunks and add up field due to each chunk (superposition principle).



- If small enough, each chunk obeys Coulomb's law

$$d\mathbf{E} = \frac{k \cdot dQ}{r^2} \hat{\mathbf{r}}$$



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Continuous charge distributions, contd

- **Discussion: Continuous charges, contd**
 - Total electric field is sum of all contributions

$$\mathbf{E} = \int d\mathbf{E}.$$

- Method can be difficult but is guaranteed to work. Next class will show easier method that works in special cases.

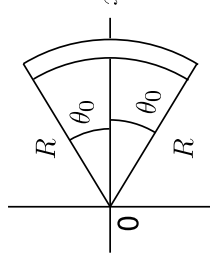


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Continuous charge distributions, contd

- **Example: Pr. 49**
 - A thin rod bent into the shape of an arc of a circle of radius R carries a uniform charge per unit length λ . The arc subtends a total angle $2\theta_0$, symmetric about the x axis, as shown below. Determine the electric field \mathbf{E} at the origin 0.



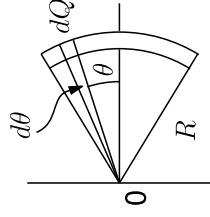
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Continuous charge distributions, contd

• Solution: Pr. 49

- We need to divide this continuous charge distribution into discrete chunks of charge dQ .
- The obvious way is to take small increments of the angle $d\theta$.



- Then the chunk has charge $dQ = \lambda R d\theta$.



Continuous charge distributions, contd

• Solution: Pr. 49, contd

$$\begin{aligned} \mathbf{E} &= \frac{k}{R^2} \int_{-\theta_0}^{+\theta_0} \lambda R d\theta (-\cos\theta \hat{\mathbf{i}} - \sin\theta \hat{\mathbf{j}}) \\ &= -\frac{k\lambda}{R} \left[\hat{\mathbf{i}} \int_{-\theta_0}^{+\theta_0} \cos\theta d\theta + \hat{\mathbf{j}} \int_{-\theta_0}^{+\theta_0} \sin\theta d\theta \right]. \end{aligned}$$

- Notice the $\hat{\mathbf{j}}$ contributions cancel out, leaving only

$$\begin{aligned} \mathbf{E} &= -\frac{k\lambda}{R} \hat{\mathbf{i}} 2 \sin\theta_0 \\ &= -\frac{2k\lambda \sin\theta_0}{R} \hat{\mathbf{i}}. \quad \square \end{aligned}$$



Continuous charge distributions, contd

• Solution: Pr. 49, contd

- Now we can apply Coulomb's law to get the electric field due to a single chunk,

$$d\mathbf{E} = \frac{k dQ}{R^2} (-\cos\theta \hat{\mathbf{i}} - \sin\theta \hat{\mathbf{j}}).$$

- We can sum over all the chunks to get the total electric field,

$$\begin{aligned} \mathbf{E} &= \int d\mathbf{E} \\ &= \frac{k}{R^2} \int dQ (-\cos\theta \hat{\mathbf{i}} - \sin\theta \hat{\mathbf{j}}) \end{aligned}$$



End

• Practice Problems:

- Ch. 21: Q. 15, 17, 19, 21, 23
- Ch. 21: Pr. 11, 13, 15, 19, 25, 27, 29, 35, 37, 39, 41, 43, 55, 57, 71, 73, 75, 77, 79, 81, 83, 87

• Interactive Quiz: Feedback

• Tutorial Question: tut05

